

RESIT PDE 2020

1) Let $w = |\nabla u|^2 = u_x^2 + u_y^2$

Since $u \in C^1(\bar{\Omega})$, w is continuous on $\partial\Omega$.

Therefore, since $\partial\Omega$ is closed & bounded,
 $\exists M$ s.t.

$$w \leq M \text{ on } \partial\Omega$$

An explicit computation gives

$$\begin{aligned} \Delta w = 2(& u_x u_{xxx} + u_{xx}^2 + u_y u_{xxy} + u_{xy}^2 \\ & + u_x u_{xyy} + u_{xy}^2 + u_y u_{yyy} + u_{yy}^2) \end{aligned}$$

Since u is harmonic:

$$\begin{aligned} u_x u_{xxx} + u_x u_{xyy} &= u_x (u_{xx} + u_{yy})_x = 0 \\ u_y u_{xyy} + u_y u_{yyy} &= u_y (u_{xx} + u_{yy})_y = 0 \end{aligned}$$

thus

$$\Delta w = 2(u_{xx}^2 + 2u_{xy}^2 + u_{yy}^2) \geq 0 \text{ in } \Omega$$

i.e. w is sub-harmonic in Ω , implying
 $|\nabla u|^2 = w \leq M$ in $\bar{\Omega}$.

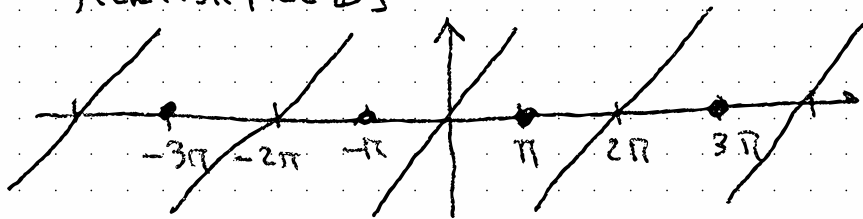
Remark this can be solved also via FTC
or mean value theorem, along w.
maximum principle.

(2) (a) As $f(x)$ is odd, all the cosine FS coefficients will be 0. We need only to find the sine Fourier series coefficients

$$\begin{aligned}
 b_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin nx \, dx \\
 &= \frac{1}{\pi} \int_{-\pi}^{\pi} x \sin nx \, dx \\
 &= \frac{1}{\pi} \left(-\frac{\cos nx}{n} \Big|_{-\pi}^{\pi} - \int_{-\pi}^{\pi} \left(-\frac{\cos nx}{n} \right) dx \right) \\
 &= \frac{2}{n\pi} \left(\frac{\sin n\pi}{n} - \pi \frac{\cos n\pi}{n} \right) \\
 &= \frac{2}{n} (-1)^{n+1}
 \end{aligned}$$

$$\Rightarrow f(x) = \sum_{n=1}^{\infty} \frac{2}{n} (-1)^{n+1} \sin nx$$

(b) The series converges to the odd periodic extension of $f(x) = x$ in $\mathbb{R} \setminus \{(2k+1)\pi \mid k \in \mathbb{Z}\}$, and $f(x) = 0$ in $\{(2k+1)\pi \mid k \in \mathbb{Z}\}$



(2) (i) $f(x) = x$ is an odd function.

Its periodic extension to \mathbb{R} , say f_{ext} , is discontinuous at the points $(2k+1)\pi = x$
 ∞

$$\lim_{x \rightarrow (2k+1)\pi^+} f_{\text{ext}}(x) = -\pi$$

$$\lim_{x \rightarrow (2k+1)\pi^-} f_{\text{ext}}(x) = \pi$$

Therefore the Fourier series converges pointwise at $x = \pm\pi$ to $\frac{1}{2}(f(\pi+\pi) + f(\pi-\pi)) = 0$.

For $x \in (-\pi, \pi)$, $f_{\text{ext}}(x)$ is continuous so the Fourier series converges pointwise to x .

$$\text{Finally } \|f\|^2 = \int_{-\pi}^{\pi} f(x)^2 dx < +\infty$$

\Rightarrow The Fourier sine series converges in L^2 sense.

Uniform convergence is not possible due to jump discontinuities and the Gibbs phenomenon.

(3) (a) The free space Green's function is the solution G_0 of

$$-G_0'' + \omega^2 G_0 = \delta(x-\xi) \quad x, \xi \in \mathbb{R}$$

(b) Assume G_0 to be known, then the superposition principle implies that

$$u(x) = \int_{\mathbb{R}} h(\xi) G_0(x; \xi) d\xi$$

(c) By applying \mathcal{F} to the equation, we get

$$\begin{aligned} \hat{u}(k)(k^2 + \omega^2) &= \hat{h}(k) \\ \Rightarrow \hat{u}(k) &= \frac{\hat{h}(k)}{k^2 + \omega^2} = \hat{h}(k) \cdot \mathcal{F}\left(\frac{1}{\sqrt{2}} \frac{e^{-\omega|x|}}{\omega}\right) \\ &= \sqrt{2\pi} \hat{h}(k) \mathcal{F}\left(\frac{1}{2} \frac{e^{-\omega|x|}}{\omega}\right) \end{aligned}$$

By the convolution formula

$$u(x) = h * f = \int_{\mathbb{R}} h(\xi) \frac{e^{-\omega|x-\xi|}}{2\omega} d\xi$$

(4) (a) Solve the eqn for the characteristic curves

$$\frac{dy}{dx} = -3x^2$$

to get $y = -x^3 + C$, where C is the constant of integration.

Solving for C we have

$$C = y + x^3$$

and thus the general solution is of the form

$$u(x, y) = f(C) = f(x^3 + y)$$

where f is an arbitrary function on \mathbb{R} .

(b) By the previous point

$$u(0, y) = f(y) = -y^2$$

i.e. $f(z) = -z^2$ and

$$u(x, y) = -(x^3 + y)^2.$$

(5) (a) let $u(r, \theta) = R(r)\theta(\theta)$, substituting in

$$\Delta u = u_{rr} + \frac{1}{r} u_r + \frac{1}{r^2} u_{\theta\theta} = 0$$

$$\text{we get } R''\theta + \frac{1}{r} R'\theta + \frac{1}{r^2} R\theta'' = 0$$

Dividing by θ and R we get

$$r^2 \frac{R''}{R} + r \frac{R'}{R} + \frac{\theta''}{\theta} = 0$$

which separates into

$$r^2 \frac{R''}{R} + r \frac{R'}{R} = -\frac{\theta''}{\theta} = \lambda,$$

i.e. the equations

$$\begin{cases} r^2 R'' + r R' - \lambda R = 0 \\ \theta'' + \lambda \theta = 0 \end{cases}$$

(b) Boundary conditions on θ imply

$$\theta(0) = \theta(\pi) = 0$$

• For $0 < \lambda = k^2$ we have solutions

$$\theta(\theta) = A \cos(k\theta) + B \sin(k\theta)$$

$$\text{From } \theta(0) = A = 0, \theta(\pi) = A \cos(k\pi) + B \sin(k\pi) = 0$$

we get $A = 0, B \sin(k\pi) = 0$, i.e.

$$k_n = n, \quad n = 1, 2, 3, \dots$$

- For $\lambda = 0$ we have

$$\theta(\theta) = A\theta + B$$

so $\theta(0) = B = 0$ and $\theta(\pi) = A\pi + B = 0$

which gives the trivial solution

$A = B = 0$, which is not a valid solution.

- For $0 > \lambda = -\gamma^2$ we have

$$\theta(\theta) = Ae^{\gamma\theta} + Be^{-\gamma\theta}$$

\Rightarrow The boundary conditions imply

$$\theta(0) = A + B = 0$$

$$\theta(\pi) = Ae^{\gamma\pi} + Be^{-\gamma\pi} = 0$$

$$\Rightarrow B = -A, A(e^{\gamma\pi} - e^{-\gamma\pi}) = 0$$

If $\gamma = 0 \Rightarrow \lambda = 0$ but by assumption $\lambda < 0$

If $\gamma \neq 0 \Rightarrow A = B = 0$ which is not a valid solution

In summary, the eigenvalues are

$$\lambda_n = k_n^2 = -n^2$$

with eigenfunctions

$$\theta_n(\theta) = \sin(n\theta)$$

(c) Now that λ is known, the radial solutions are found with the ansatz $R = r^\alpha$

which gives $\alpha = \pm n \Rightarrow R(r) = A_n r^n + B_n r^{-n}$

To be well defined at $r = 0$, we need to

require $B_n = 0$

(d) Putting all the results together, we get

$$u(r, \theta) = \sum_{n \geq 1} A_n r^n \sin(n\theta)$$

For $r = a$ we have

$$u(a, \theta) = \sum_{n \geq 1} A_n a^n \sin(n\theta) = h(\theta)$$

$$\Rightarrow A_n = \frac{2}{\pi a^n} \int_0^\pi h(\theta) \sin(n\theta) d\theta$$

In our case

$$\sum_{n \geq 1} A_n a^n \sin(n\theta) = \sin \theta$$

$$\Rightarrow A_1 = 1/a, \quad A_n = 0 \text{ for } n > 1$$

which means

$$u(r, \theta) = \frac{r}{a} \sin \theta.$$